

TOWARDS PRECISION MEASUREMENTS OF



COHERENT SCATTER CROSS-SECTIONS

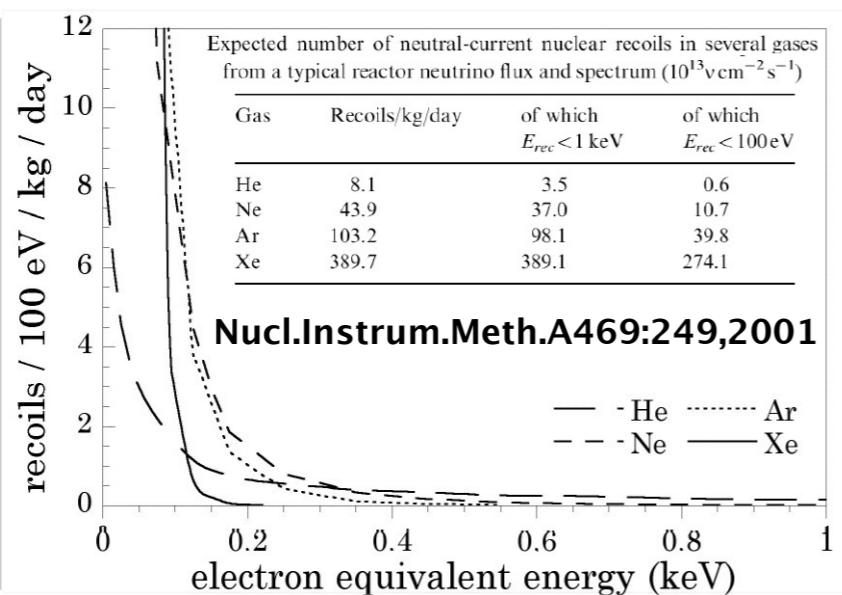
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WHAT DO WE MEAN BY “PRECISION”

A one-page tutorial on coherent ν -N scattering

- Uncontroversial Standard Model process
 - Large enhancement in cross-section for $E_\nu <$ few tens of MeV ($\sigma \propto N^2$, possible only for neutral current)
 - However, not yet measured... detector technology has been missing.
- Detector mass might be at least ~ 1 kg (reactor experiment) + recoil energy threshold $\ll 1$ keV (low-E recoils lose only 10-20% to ionization or scintillation)
- Cryogenic bolometers and other methods proposed, no successful implementation yet

ca. 2002



Fundamental physics:

- Largest σ_ν in SN dynamics: should be measured to validate models (J.R. Wilson, PRL 32 (74) 849)
- A large detector can measure total E and T of SN $\nu_\mu - \nu_\tau \Rightarrow$ determination of ν oscillation pattern and mass of ν star (J.F.Beaum, W.M.Ford & P.Vogel, PRD 66(02)03301)
- Coherent σ same for all known ν ... oscillations observed in a coherent detector \Rightarrow evidence for ν_{sterile} (A.Drukier & L.Stodolsky, PRD 30 (84) 2295)
- Sensitive probe of weak nuclear charge \Rightarrow test of radiative corrections due to new physics above weak scale (L.M.Krauss, PLB 269, 407)
- More sensitive to NSI and new neutral bosons than ν factories. Also effective ν charge ratio (J. Barranco et al, hep-ph/0508299, hep-ph-0512029)
- σ critically depends on μ_ν : observation of SM prediction would increase sensitivity to μ_ν by an order of magnitude (A.C.Dodd et al, PLB 266 (91) 134)

Smallest detectors... “ ν technology”?

- Monitoring of nuclear reactors against illicit operation or fuel diversion: present proposals using conventional 1-ton detectors reach only $> \sim 3$ GWT reactor power
- Geological prospection, planetary tomography... the list gets much wilder.

LOW HANGING FRUIT?

- ✿ NC sterile \square search
- ✿ Coherent scattering $\mu \square$ search
- ✿ Test Q_W radiative corrections
- ✿ NSI \square cross-sections
- ✿ Light Wimps

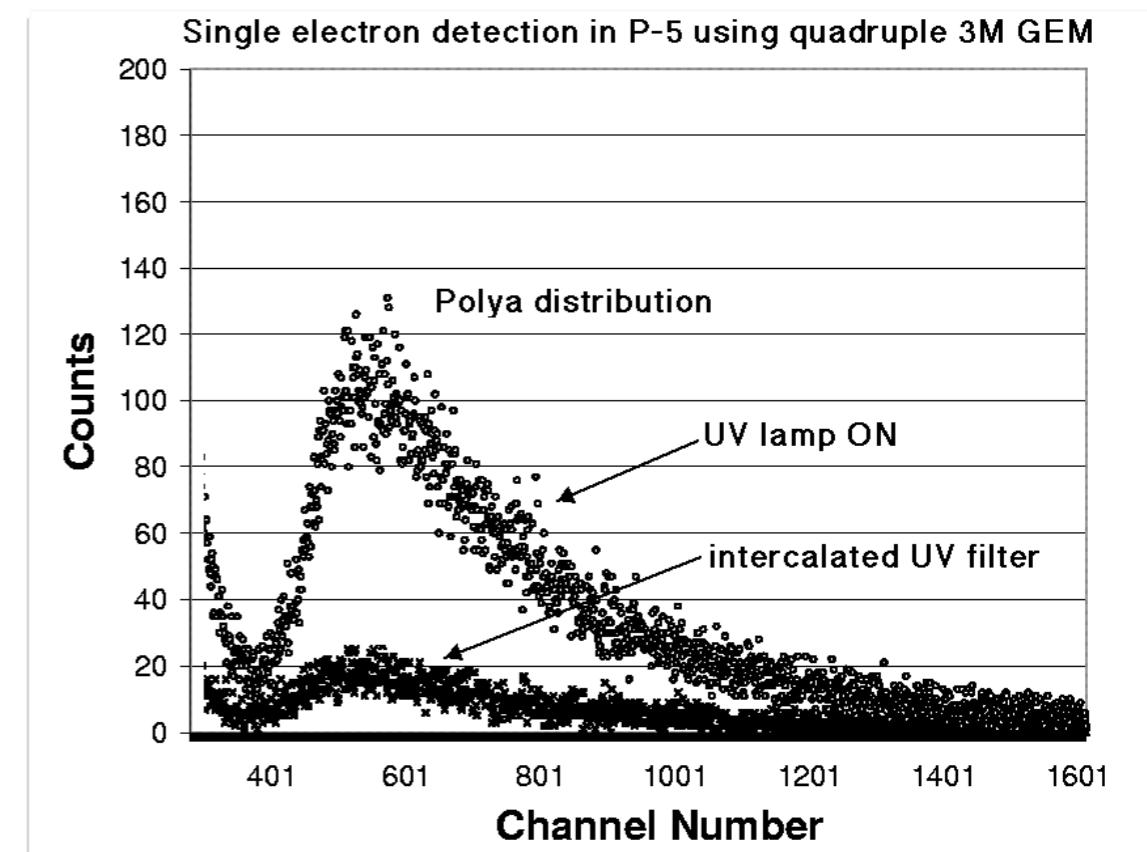
Not with scheme proposed here

AN OLD EFFORT: LOW THRESHOLD GAS DETECTORS

- Identical detector construction for all targets (controls systematics)
- Continuously re-purify
- Low thresholds: single electron sensitivity for ~1-10 Bar
- High Precision quenching factor measurement
- Target Masses ($\sim \text{kg/m}^3$)
- Standard drift gas targets: H_2 , $^{3,4}\text{He}$, $^{10,11}\text{BF}_3$, $^{12,13,14}\text{CH}_4$, C_2H_6 , C_4H_{10} ... CF_4 , $^{32,34}\text{SF}_6$, CO_2 , $^{20,22}\text{Ne}$, N_2 , $^{82,83,84,85,86}\text{Kr}$, $^{39,40}\text{Ar}$, $^{129-132,134,136}\text{Xe}$
- Low backgrounds: ^{222}Rn @EXO, n & γ @CoGeNT.

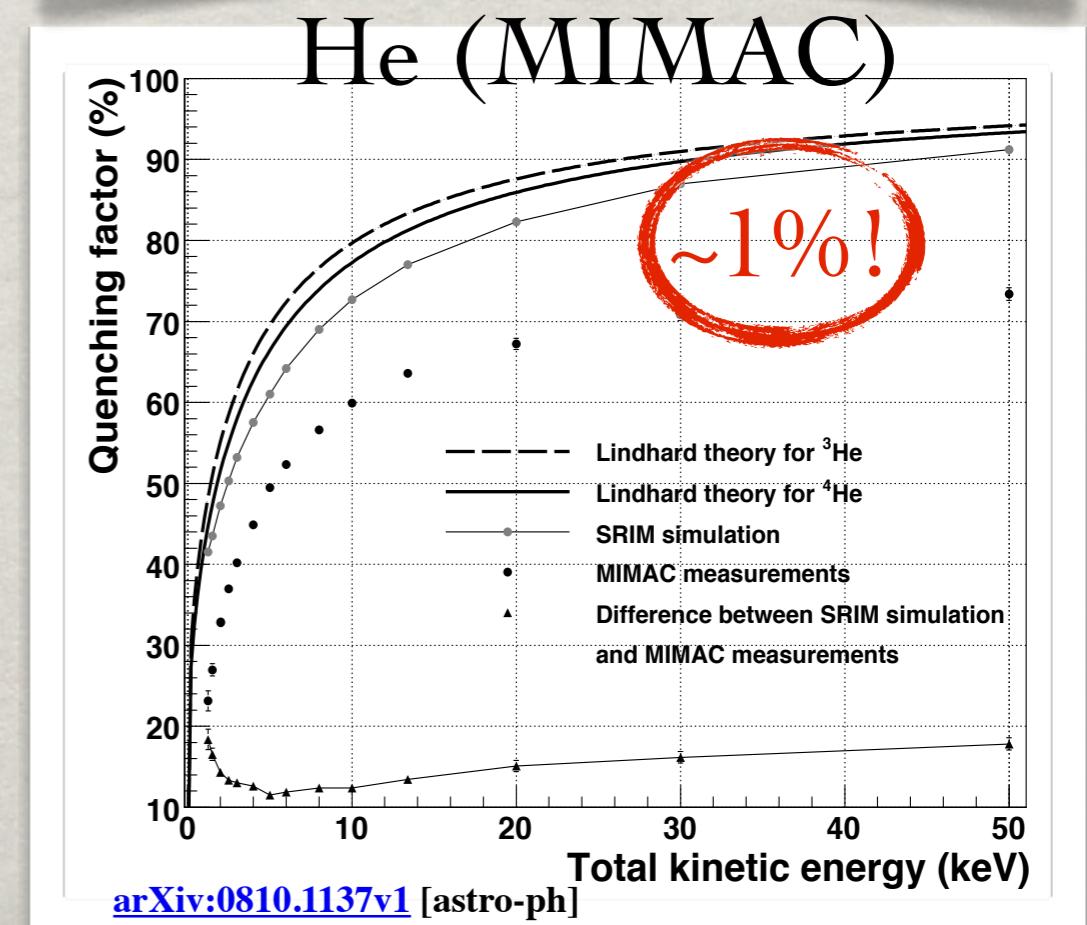
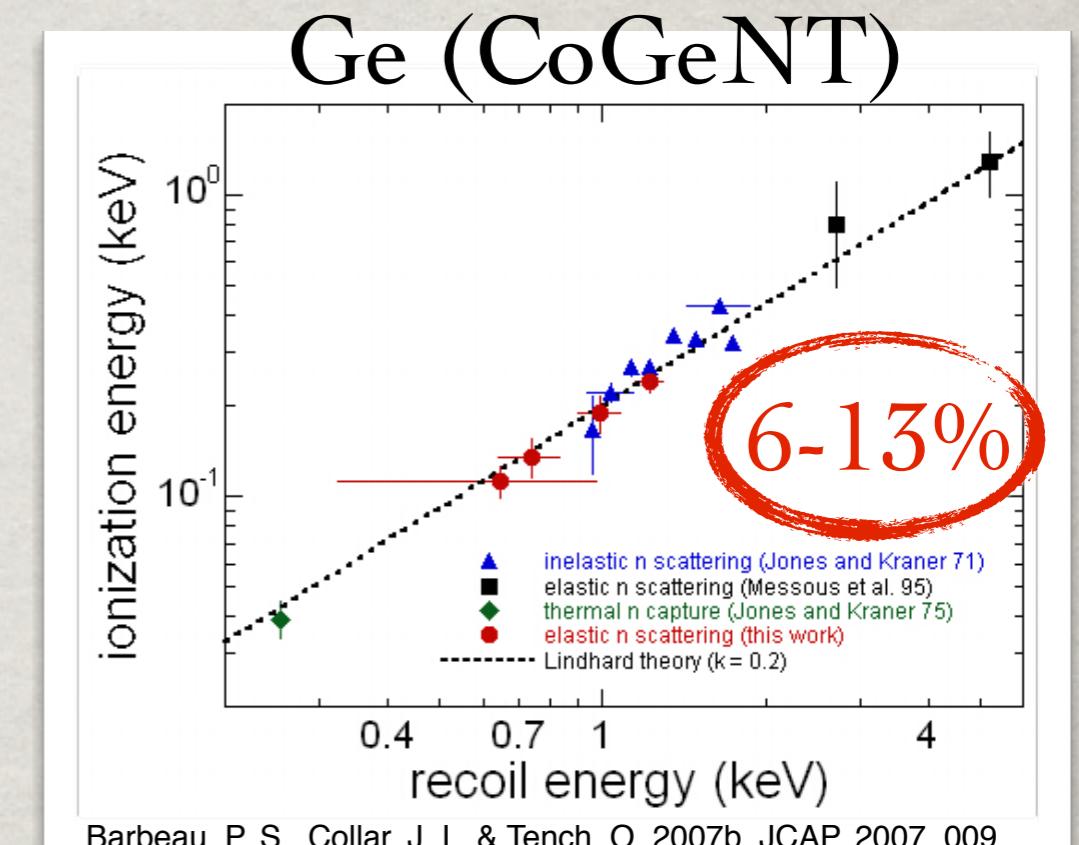


P.S. Barbeau, J.I. Collar et al., NIM A515:439– 445, 2003.



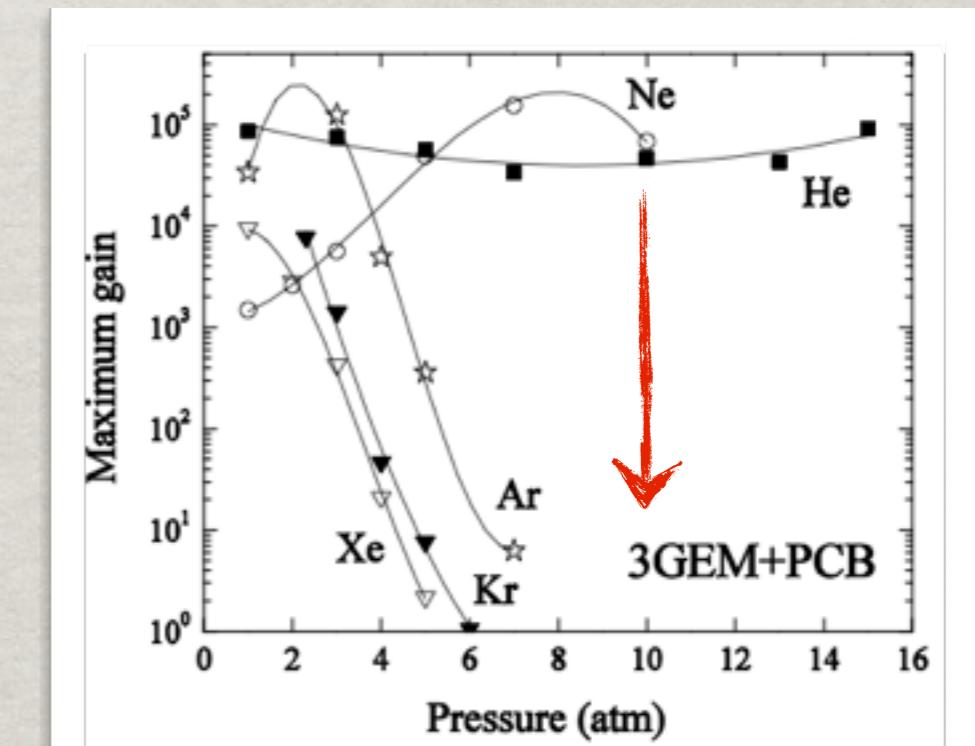
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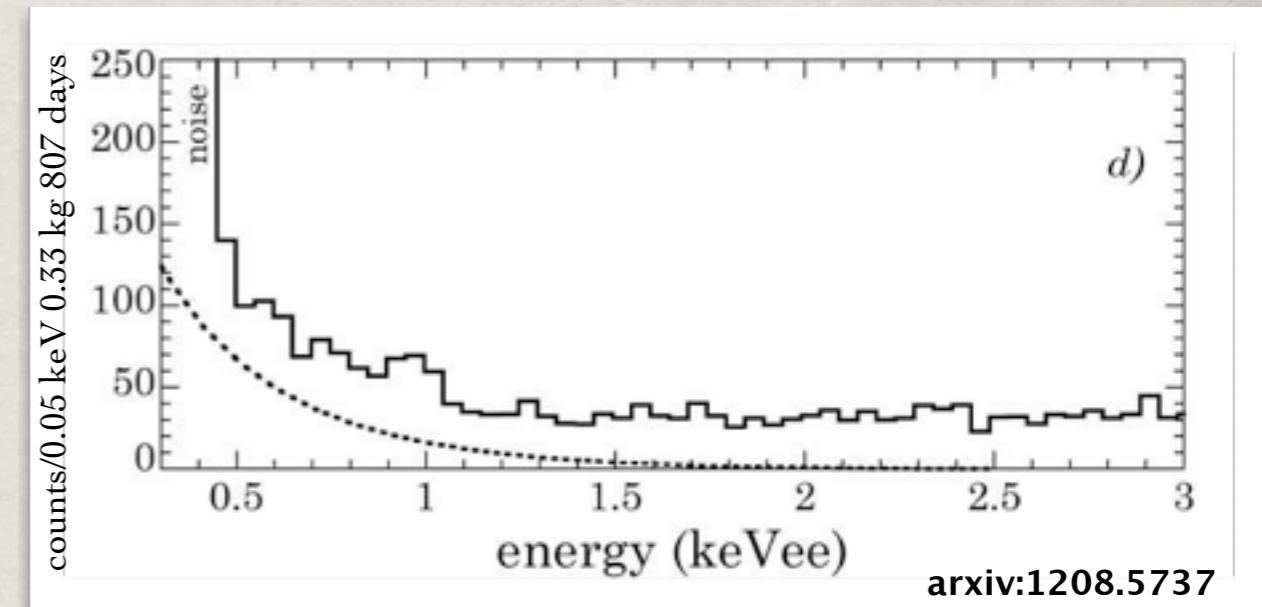
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Nucl.Instrum.Meth. A493 (2002) 8-15
Fig.8 Maximum gain of a triple GEM detector as a function of pressure in He, Ne, Ar, Kr and Xe.

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- ✿ Identical detector construction for all targets (controls systematics)
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- ✿ Low backgrounds: ²²²Rn @EXO, n & γ @CoGeNT.



²²²Rn in EXO-200 (continuously circulating):
 $4.6 \mu\text{Bq kg}^{-1} \rightarrow \sim 0.01\text{-}0.1 \text{ c keV}^{-1} \text{ kg}^{-1} \text{ d}^{-1}$

⁸⁵Kr in Xe: $< \sim 10^{-3} - 0.5 \text{ c keV}^{-1} \text{ kg}^{-1} \text{ d}^{-1}$

²²²Rn daughters on SNO NCD surfaces:
 $\sim 2 \text{ m}^{-2} \text{ d}^{-1} \rightarrow \sim 0.5 \text{ c keV}^{-1} \text{ kg}^{-1} \text{ d}^{-1}$

¹⁴C @ Borexino levels (measured in CH₄):
 $^{14}\text{C}/^{12}\text{C} < 10^{-18} \rightarrow < \sim 0.15 \text{ c keV}^{-1} \text{ kg}^{-1} \text{ d}^{-1}$

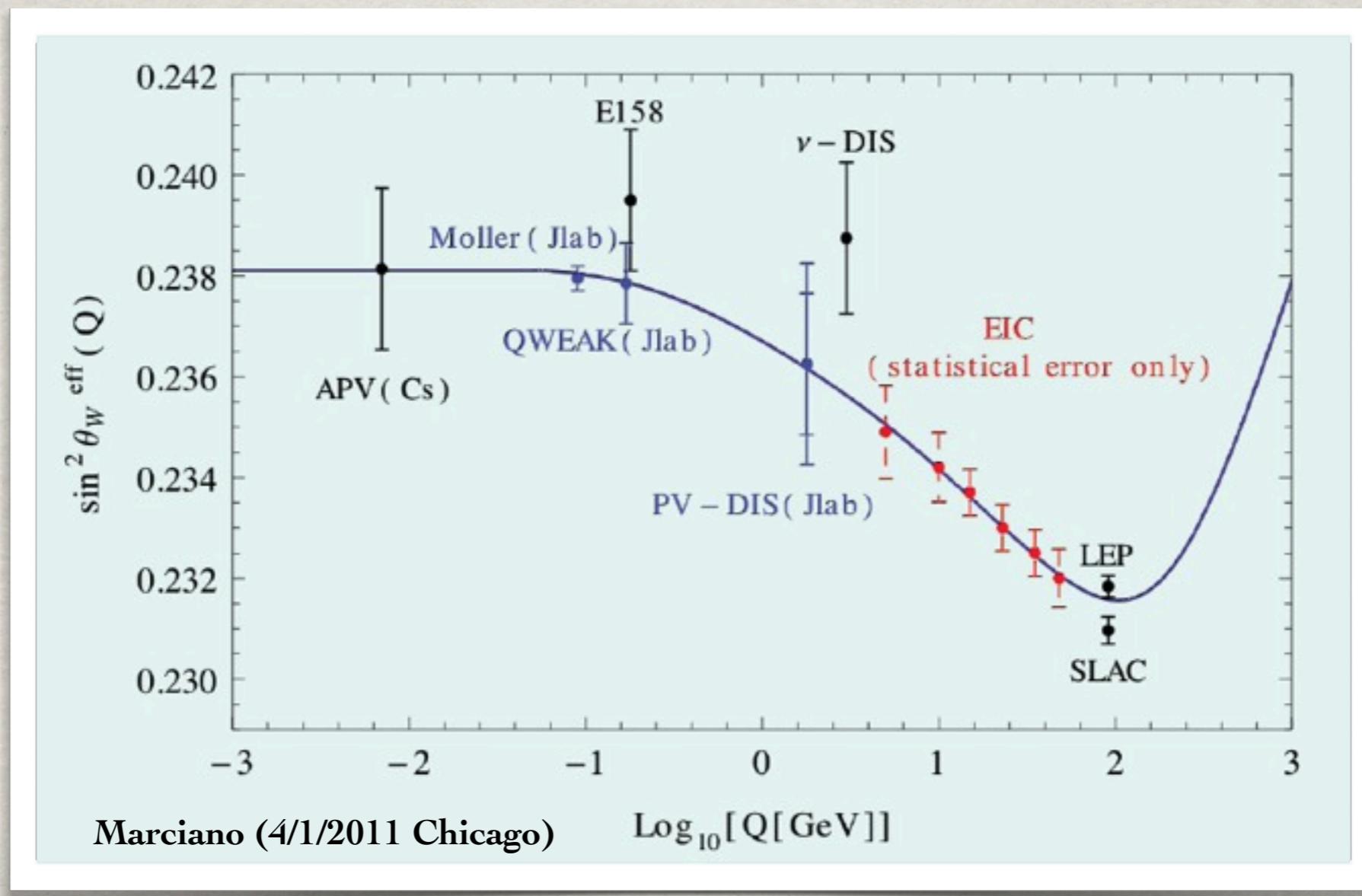
³⁹Ar in Ar @ $\sim 15 - 300 \text{ c keV}^{-1} \text{ kg}^{-1} \text{ d}^{-1}$

CoGeNT backgrounds (0.5-3 keV)
 $\underline{2.6 - 7.4 \text{ c keV}^{-1} \text{ kg}^{-1} \text{ d}^{-1}}$

WEAK NUCLEAR CHARGE

- We now know M_{top} and $M_{Higgs} \rightarrow$ uncertainties on radiative corrections are small
- Remaining hadronic uncertainties similar to those from APV experiments ($\sim 0.2\%$)
(L. M. Krauss, PLB 269, 407)

$$Q_w = N - (1 - 4\sin^2\theta_w)Z$$



WEAK NUCLEAR CHARGE

Measure Q_w with coherent $\bar{\nu}$ scattering at nuclear reactor
(SONGS $\sim 10^{13} \text{ cm}^{-2} \text{ s}^{-1}$ & 30 m.w.e)

Deviations \rightarrow new Physics

$$\frac{d\sigma}{dT}_{coh} = \frac{G_f^2 M}{2\pi} ((G_V + G_A)^2 + (G_V - G_A)^2 (1 - \frac{T}{E_\nu})^2 - (G_V^2 - G_A^2) \frac{MT}{E_\nu^2})$$

$$G_V = ((g_v^p + 2\epsilon_{ee}^{uV} + \epsilon_{ee}^{dV})Z + (g_v^n + \epsilon_{ee}^{uV} + 2\epsilon_{ee}^{dV})N)F_{nucl}^V(Q^2)$$

$$G_A = ((g_a^p + 2\epsilon_{ee}^{uA} + \epsilon_{ee}^{dA})(Z_+ - Z_-) + (g_a^n + \epsilon_{ee}^{uA} + 2\epsilon_{ee}^{dA})(N_+ - N_-))F_{nucl}^A(Q^2)$$

$$g_V^p = \rho_{\nu N}^{NC} \left(\frac{1}{2} - 2\hat{\kappa}_{\nu N} \sin^2 \theta_w \right) + 2\lambda^{uL} + 2\lambda^{uR} + \lambda^{dL} + \lambda^{dR}$$

$$g_V^n = -\frac{1}{2}\rho_{\nu N}^{NC} + \lambda^{uL} + \lambda^{uR} + 2\lambda^{dL} + 2\lambda^{dR}$$

+ axial vector factors which have more theoretical uncertainty (strong quark contributions, weak magnetism term, effective neutrino charge radii)

WEAK NUCLEAR CHARGE

- 0) Use gas targets (swappable) to control fiducial volume systematics
- 1) Low q^2 @ Rx to avoid $F(Q^2)$ theoretical systematics
- 2) eliminate axial couplings along with their (larger) uncertainties

→ Choose even-even nuclei

H_2	$^{32,34}SF_6$
$^{3,4}He$	CO_2
$^{10,11}BF_3$	$^{20,22}Ne$
$^{12,13,14}CH_4$	N_2
C_2H_6	$^{82,83,84,85,86}Kr$
$C_4H_{10} \dots$	$^{39,40}Ar$
CF_4	$^{129-132,134,136}Xe$

WEAK NUCLEAR CHARGE

3) Factorize out \square flux ($\sim 6\%$) & absolute rate uncertainties

→ group according $Z=N$ & $Z \neq N$ & measure ratio: $\frac{R_{Z=N}}{R_{Z \neq N}}$

$$\begin{aligned}Q_{w,{}^4He} &= 2 \times 4 \sin^2 \theta_w \\Q_{w,{}^{12}C} &= 6 \times 4 \sin^2 \theta_w \\Q_{w,{}^{16}O} &= 8 \times 4 \sin^2 \theta_w \\Q_{w,{}^{20}Ne} &= 10 \times 4 \sin^2 \theta_w\end{aligned}$$

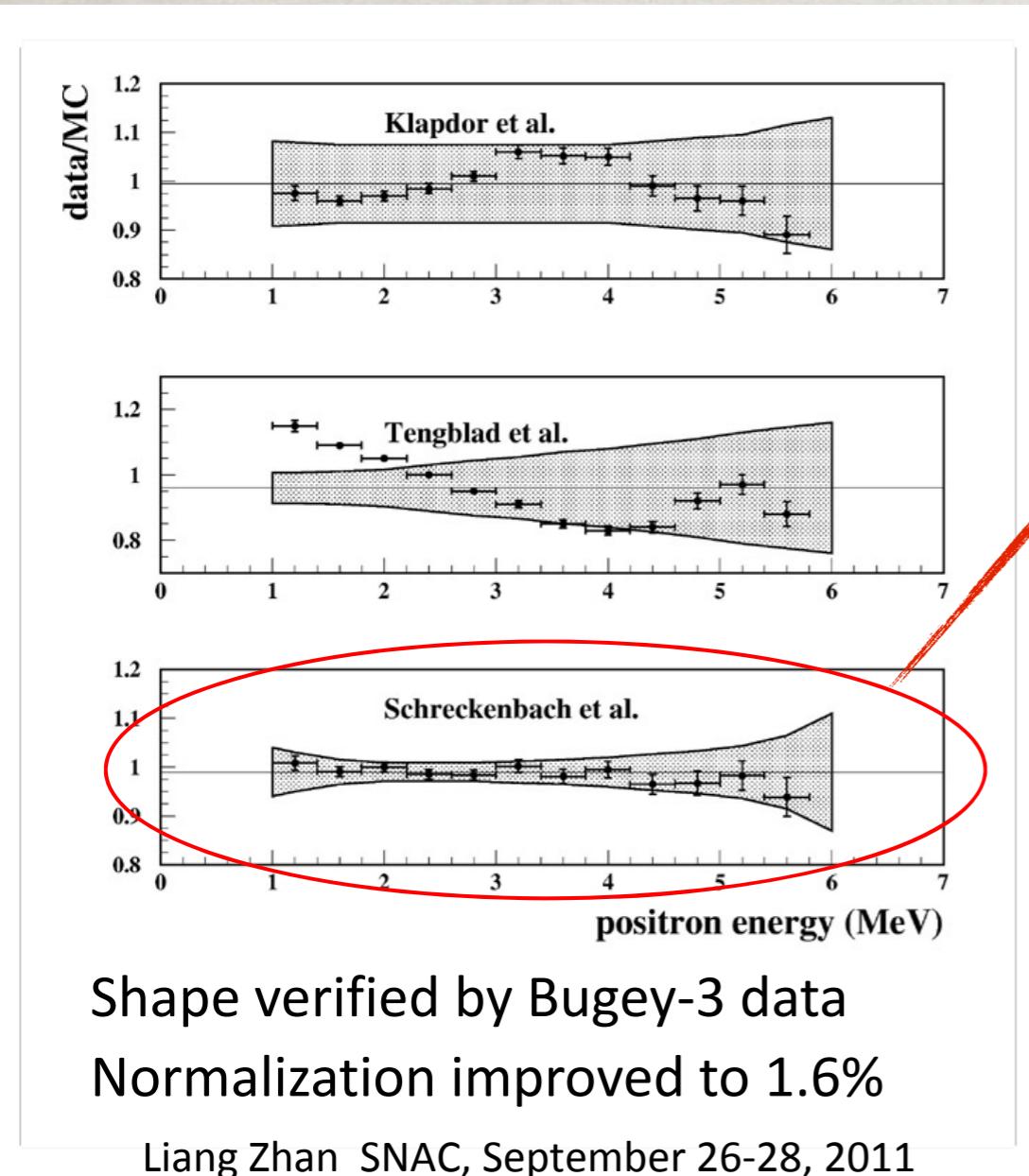
$$\begin{aligned}Q_{w,{}^{22}Ne} &= 2 + 10 \times 4 \sin^2 \theta_w \\Q_{w,{}^{40}Ar} &= 4 + 18 \times 4 \sin^2 \theta_w \\Q_{w,{}^{136}Xe} &= 28 + 54 \times 4 \sin^2 \theta_w\end{aligned}$$

$$Q_w = N - (1 - 4 \sin^2 \theta_w) Z$$

WEAK NUCLEAR CHARGE

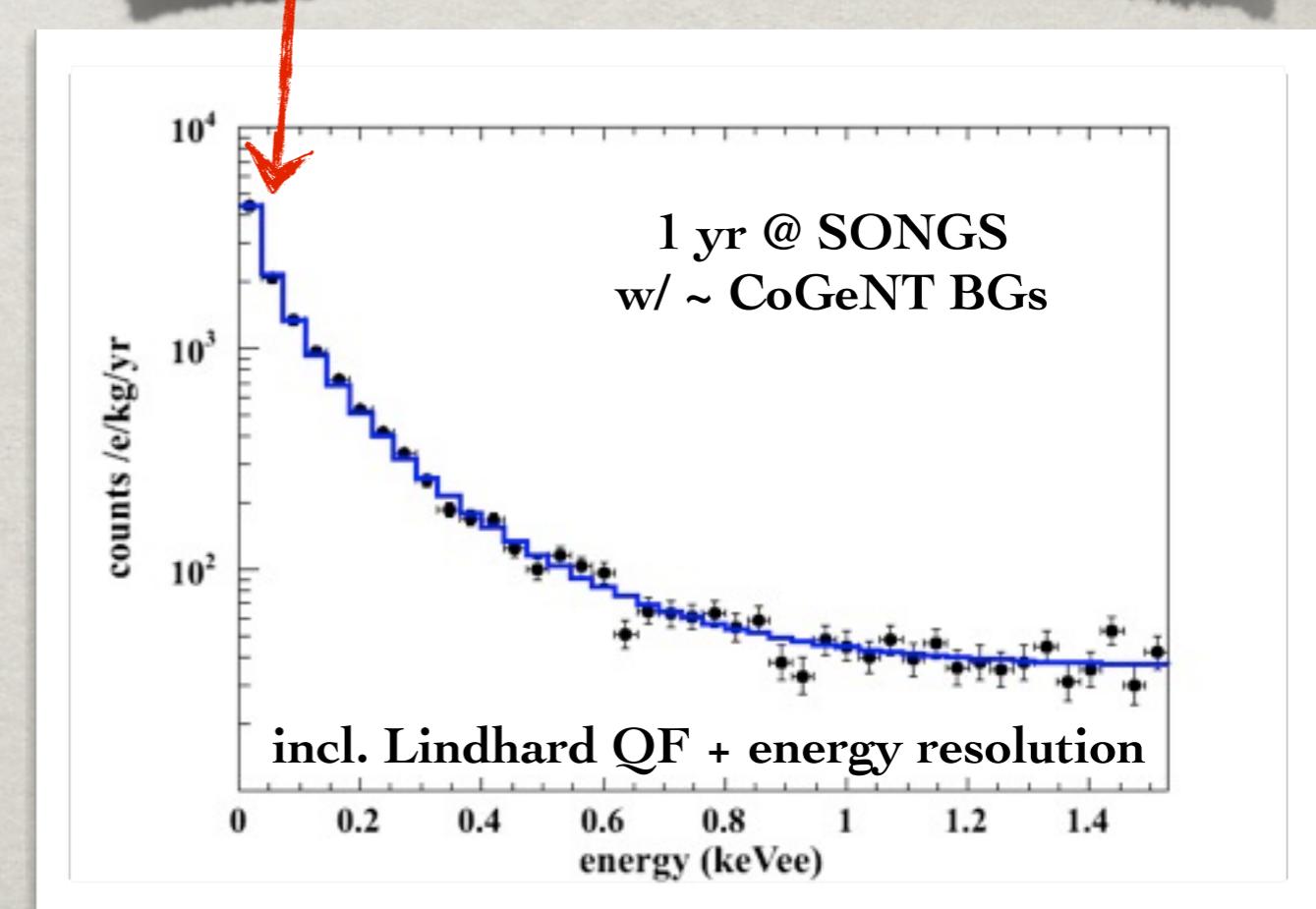
4) Use $A_1 \sim A_2$ nuclei to minimize impact of neutrino spectrum uncertainties $\rightarrow {}^{20, 22}\text{Ne}$

$$\text{Recoil energy: } T_{\max} = 2E_{\square}/M$$



Choose recoil thresholds (10% change between ${}^{20, 22}\text{Ne}$) to select same population of \square energies (spectral uncertainties factorize out)

Introduces <0.1% uncertainty due to discrete nature of the recorded signal (single e⁻'s) @ threshold.



WEAK NUCLEAR CHARGE

4.5) Using same element (Ne) eliminates atomic effects on the quenching factor

impact of 1% $Q(E_{\text{rec}})$ uncertainty & threshold

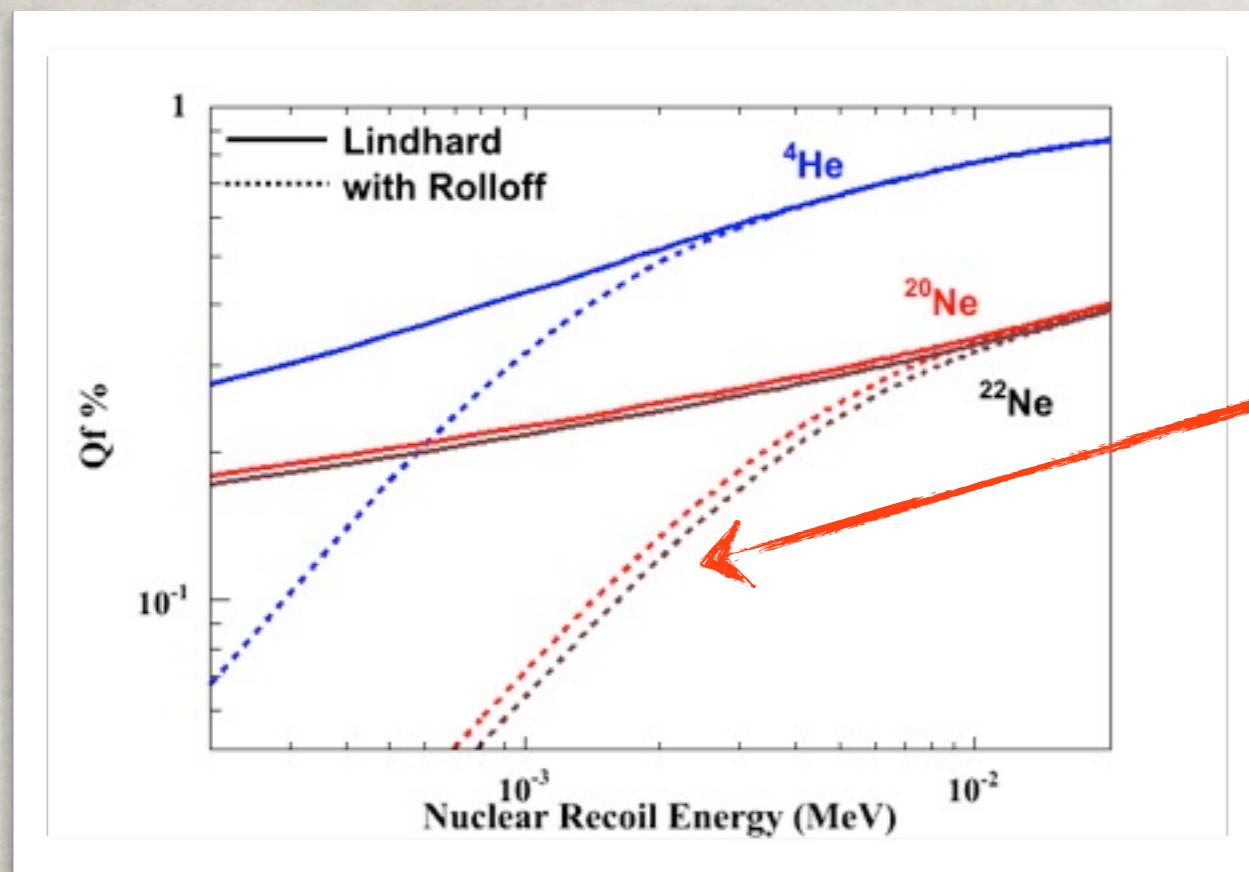
Threshold (e ⁻ 's)	Systematic impact (%)
0	0.1%
1	0.4%
2	0.6%
3	0.8%

If we measure the Ratio of the Quenching Factor in $^{20,22}\text{Ne}$ to $\sim 1\%$, then the systematics are manageable

From Lindhard, this kinematic change comes in as:

$$f_n \sim \frac{1}{A^{\frac{1}{2}} + 1} \times \left(1 - e^{\frac{-E_r}{E_t}}\right), \quad E_t \sim A$$

Should be able to predict the difference; but should still measure that it is non-zero. Can test ratio with other targets $^{3,4}\text{He}$.



WEAK NUCLEAR CHARGE

Statistical uncertainty from backgrounds dominate.

→ Need Rx-off time

Run for 4.5 cycles at SONGS. 1 cycle = 18 mo. On, 1 mo. Off (*When they are operating normally*)

Operate in both Tendon Galleries to maximize Rx off time.

→ 2 x 20 kg detectors at ~ 1-10 Bar

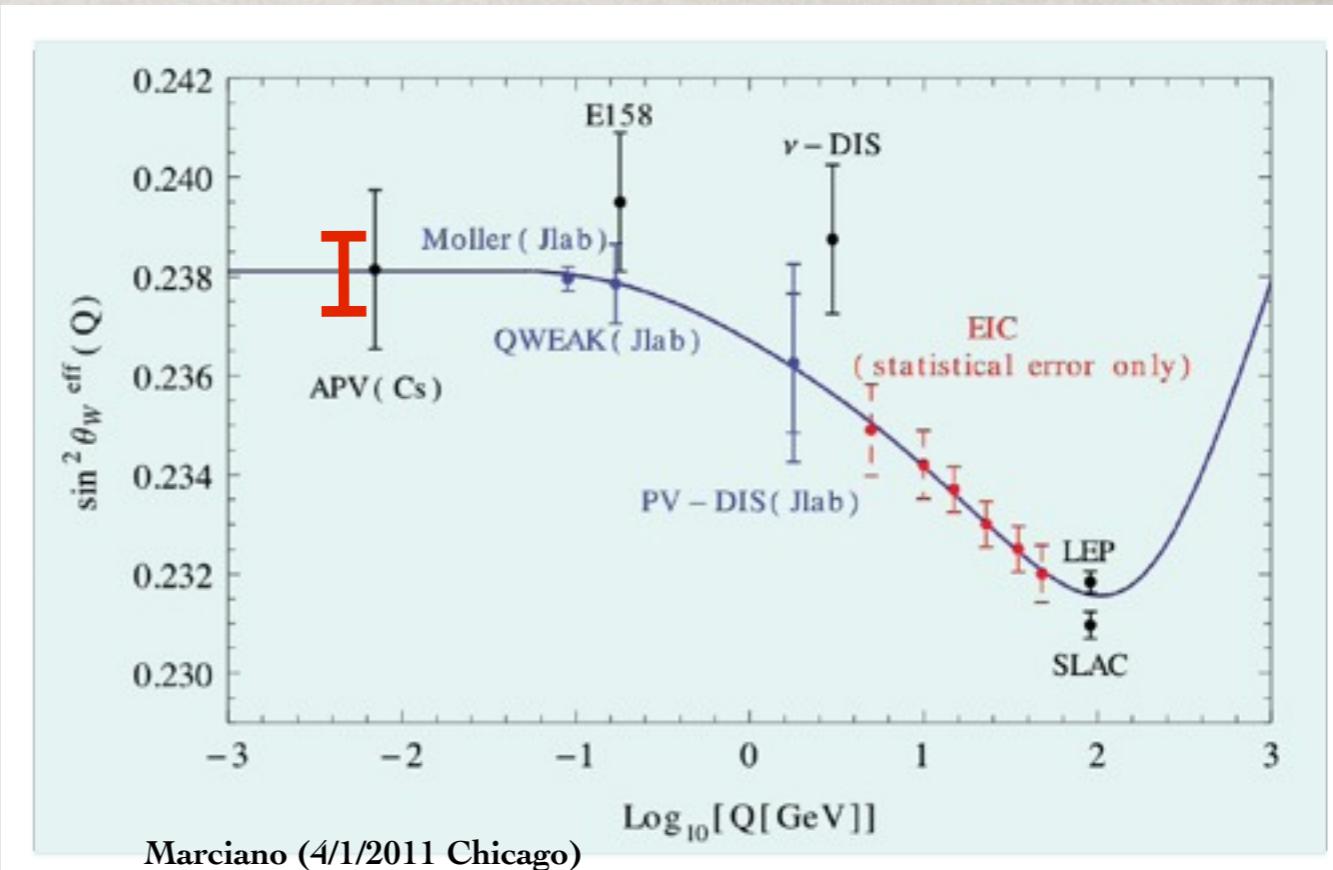
Result → uncertainties on $\sin^2 \theta_w$:
 $\pm 0.22\%$ (stat.) $\pm [0.1-0.4]\%$ (sys.) $\pm <0.2$ (th.)

Gives us another neutrino test, at lower Q .

Ignoring radiative corrections

$$R\left(\frac{^{22}Ne}{^{20}Ne}\right) = \frac{(2 + 10 \times \sin^2 \theta_w)^2}{(10 \times \sin^2 \theta_w)^2}$$

$$\sigma(\sin^2 \theta_w) = 0.57 \times \sigma R$$



(NON-UNIVERSAL) NSI SEARCH

- Essentially, the same game as the before

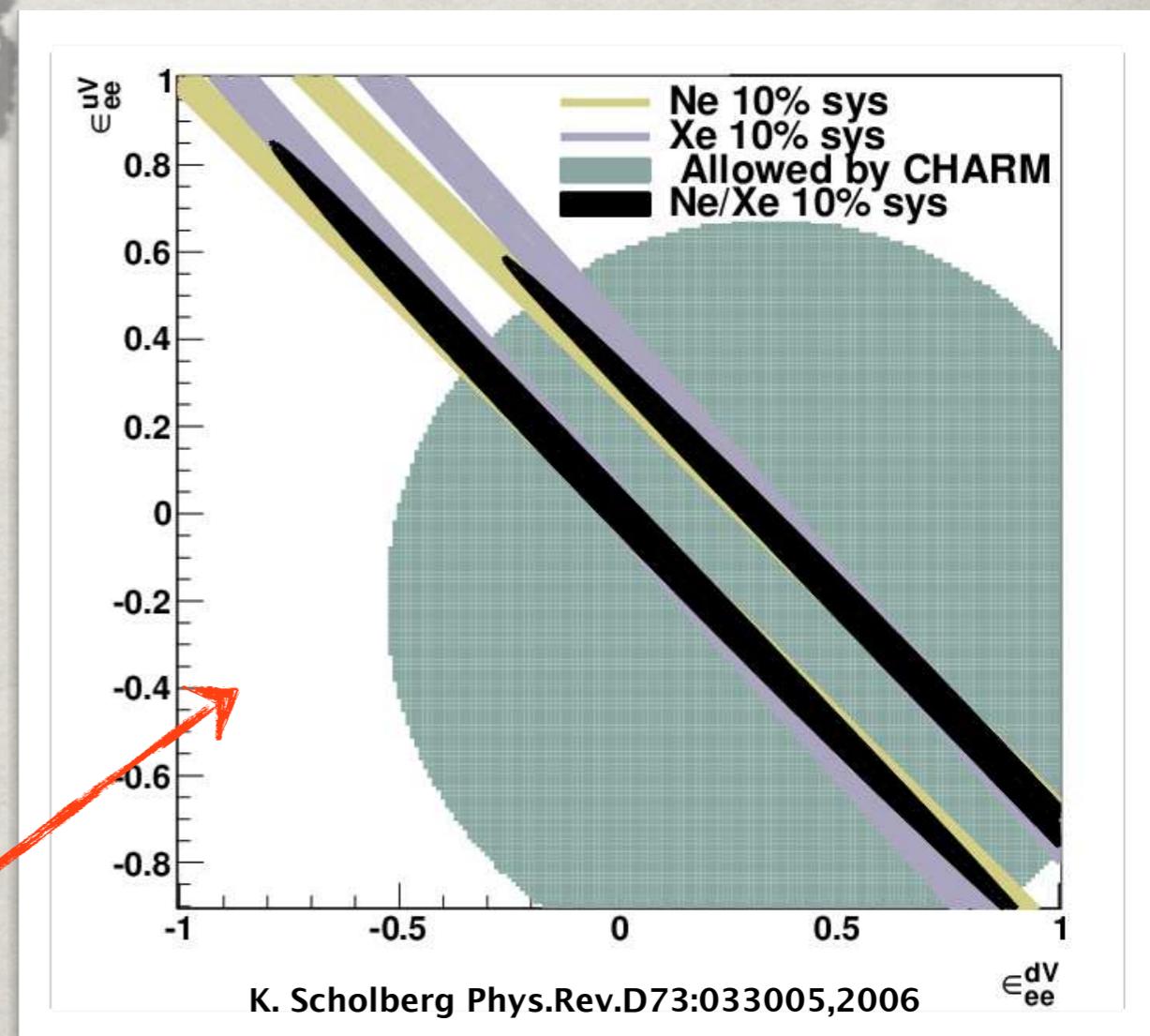
$$\frac{d\sigma}{dT}_{coh} = \frac{G_f^2 M}{2\pi} = G_V^2 \left(1 + \left(1 - \frac{T}{E_\nu}\right)^2 - \frac{MT}{E_\nu}\right)$$

$$G_V = ((g_v^p + 2\epsilon_{ee}^{uV} + \epsilon_{ee}^{dV})Z + (g_v^n + \epsilon_{ee}^{uV} + 2\epsilon_{ee}^{dV})N)F_{nucl}^V(Q^2)$$

Including radiative corrections, and earlier stat. & sys. uncertainties, the ratio for ${}^{20}\text{Ne}/{}^{22}\text{Ne}$ gives:

$$1.0345 \pm 0.0202 = \frac{-0.512 + \epsilon_{ee}^{uV} + 2\epsilon_{ee}^{dV}}{-0.495 + 3\epsilon_{ee}^{uV} + 3\epsilon_{ee}^{dV}}$$

Makes for interesting constraints here
(not yet drawn)



NSI SEARCH

While we are at it, lets not forget that this scheme is information rich

- H₂, CH₄ (very distinctive spectrum) and CF₄ (unpaired protons)
- ³He (unpaired neutron), D₂ (unpaired neutron and proton) □ Can't decide which is crazier
- BF₃ (unpaired neutron and proton).
- Varying weak magnetism effect.

$$\frac{d\sigma}{dT}_{coh} = \frac{G_f^2 M}{2\pi} ((G_V + G_A)^2 + (G_V - G_A)^2 (1 - \frac{T}{E_\nu})^2 - (G_V^2 - G_A^2) \frac{MT}{E_\nu^2})$$

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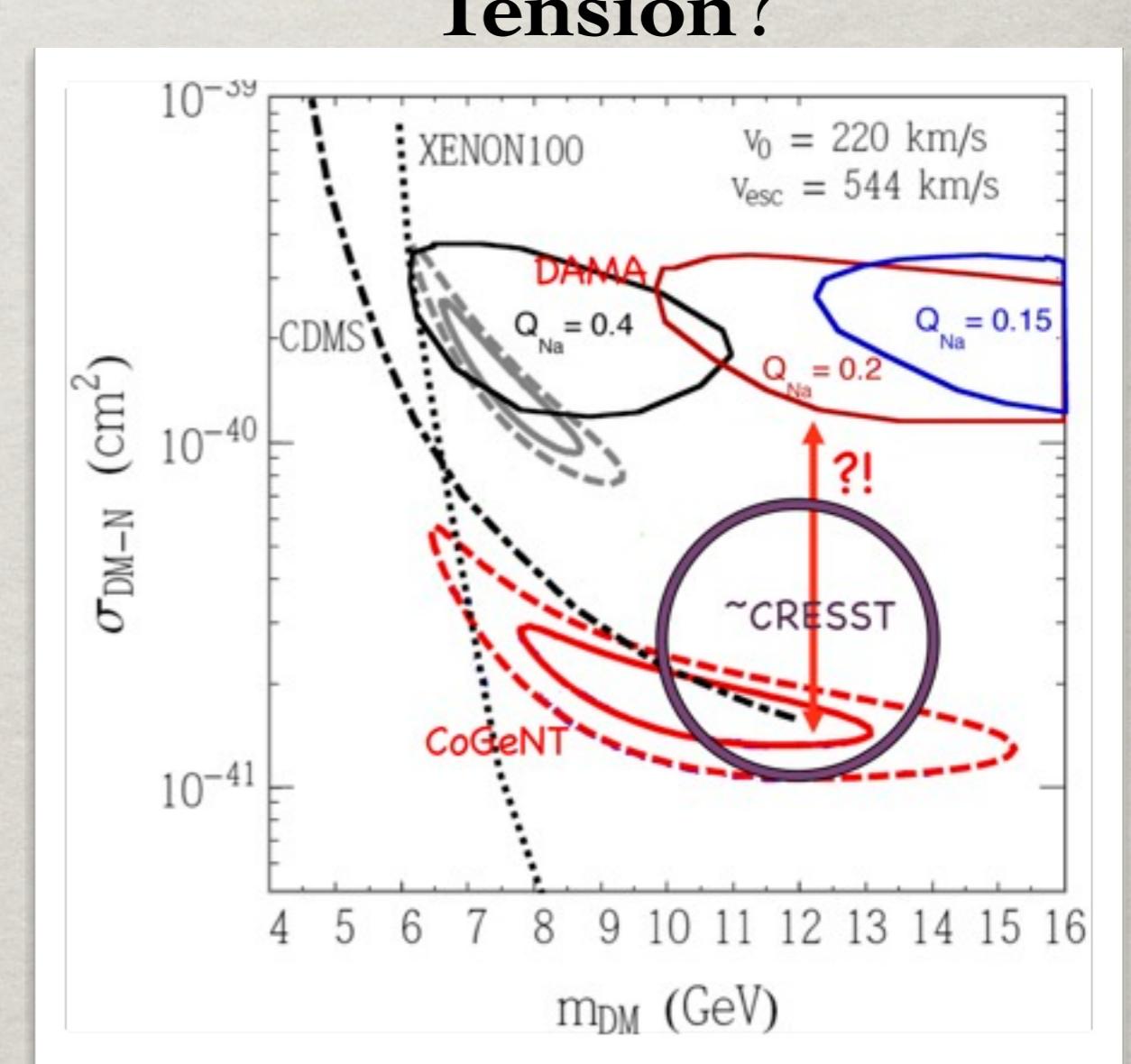
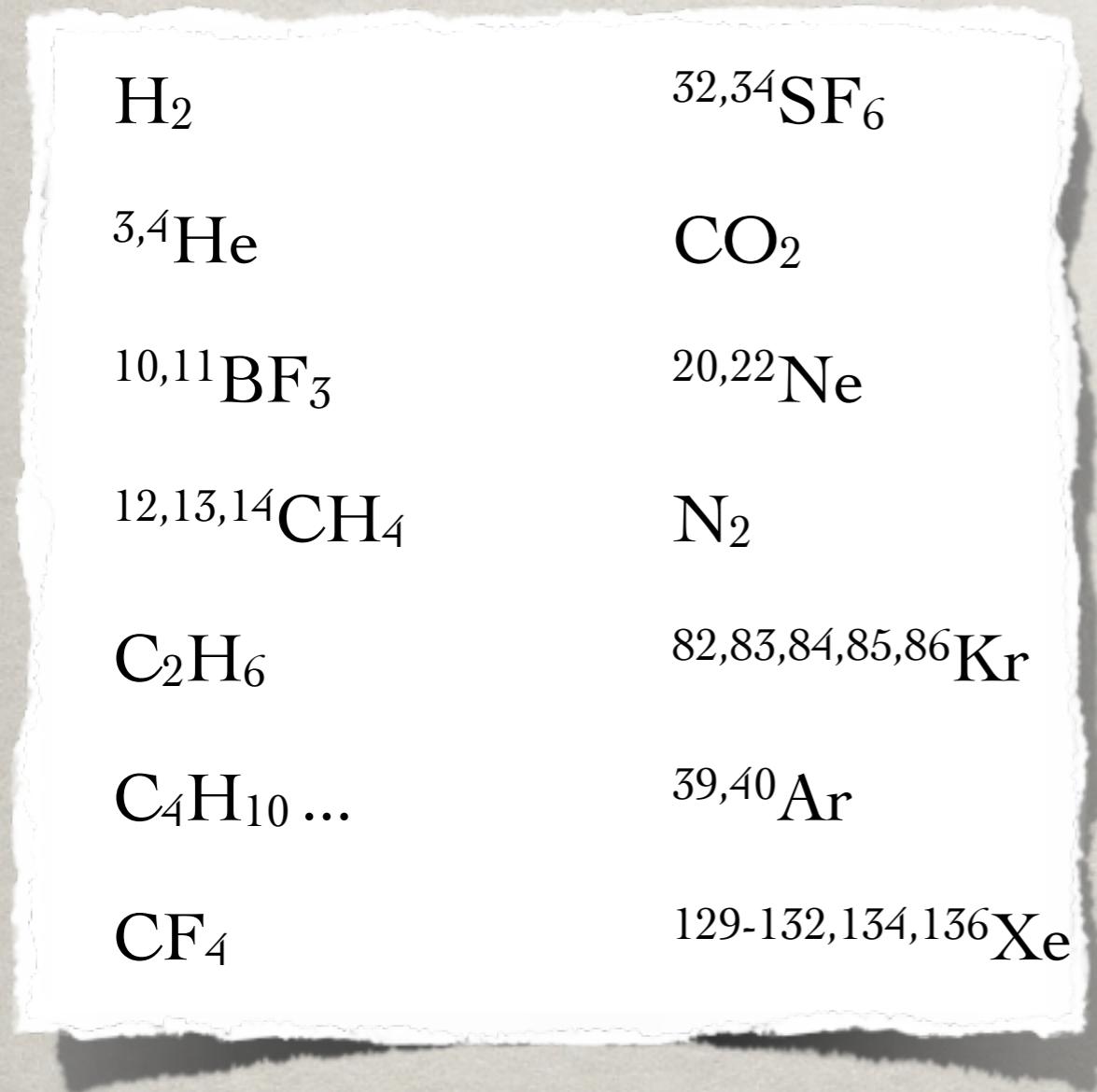
$$g_V^n = -\frac{1}{2}\rho_{\nu N}^{NC} + \lambda^{uL} + \lambda^{uR} + 2\lambda^{dL} + 2\lambda^{dR} \quad \text{+ pesky axial couplings}$$

Light WIMPS

Or, what else can you use these detectors for...?

Deploy a number of similarly built detectors, with much larger variance in A.

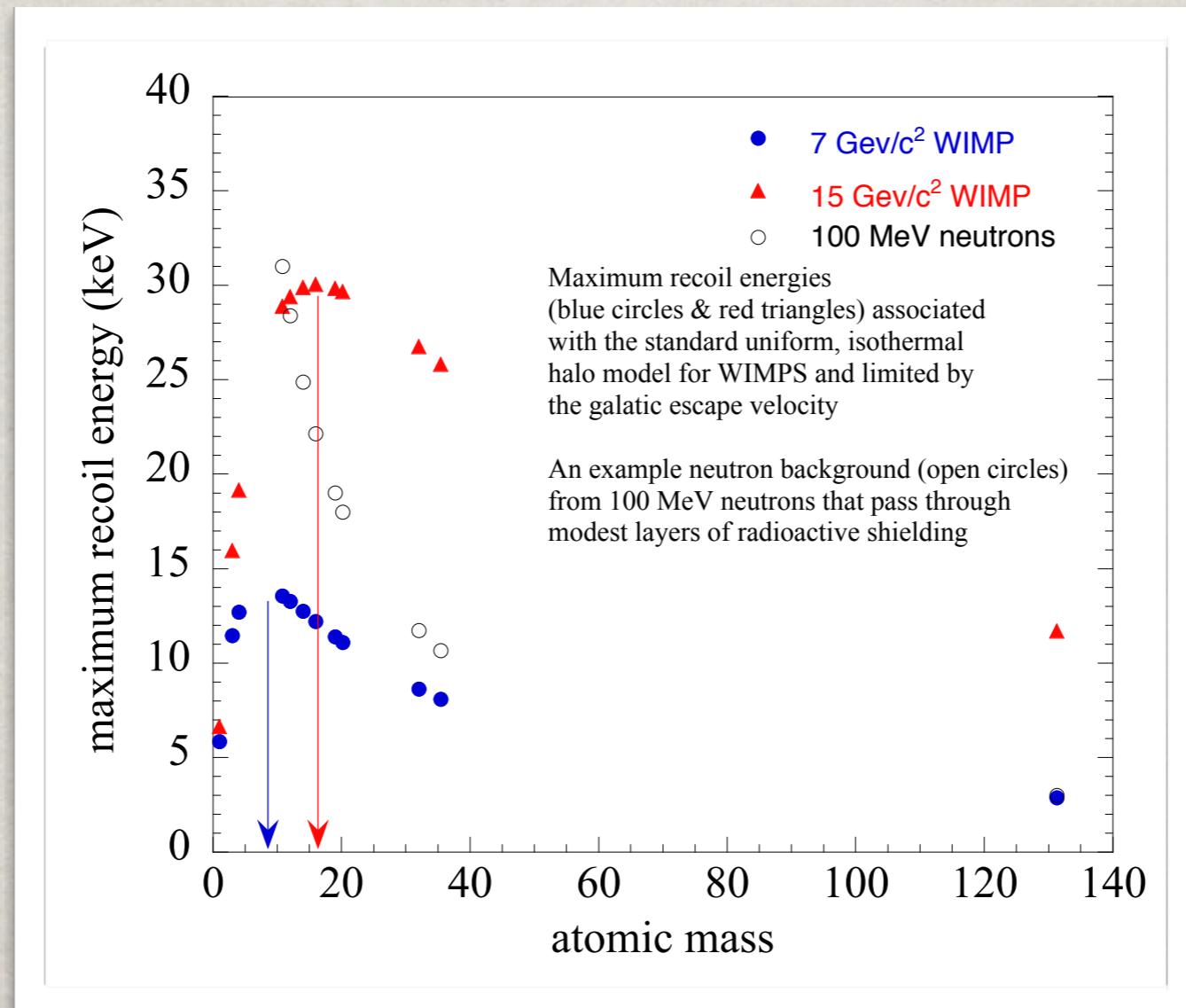
Use the kinematics of WIMP-nucleus scattering to test putative signals.



Light WIMPS

Fit characteristic energy scale of any observation versus target mass, increases the precision on putative WIMP mass
→ kinematic check against certain (neutron) background hypotheses

Amplitude of spectra indicates WIMP escape velocity
→ ascertain/factorize astrophysical systematic (Streams, etc.)



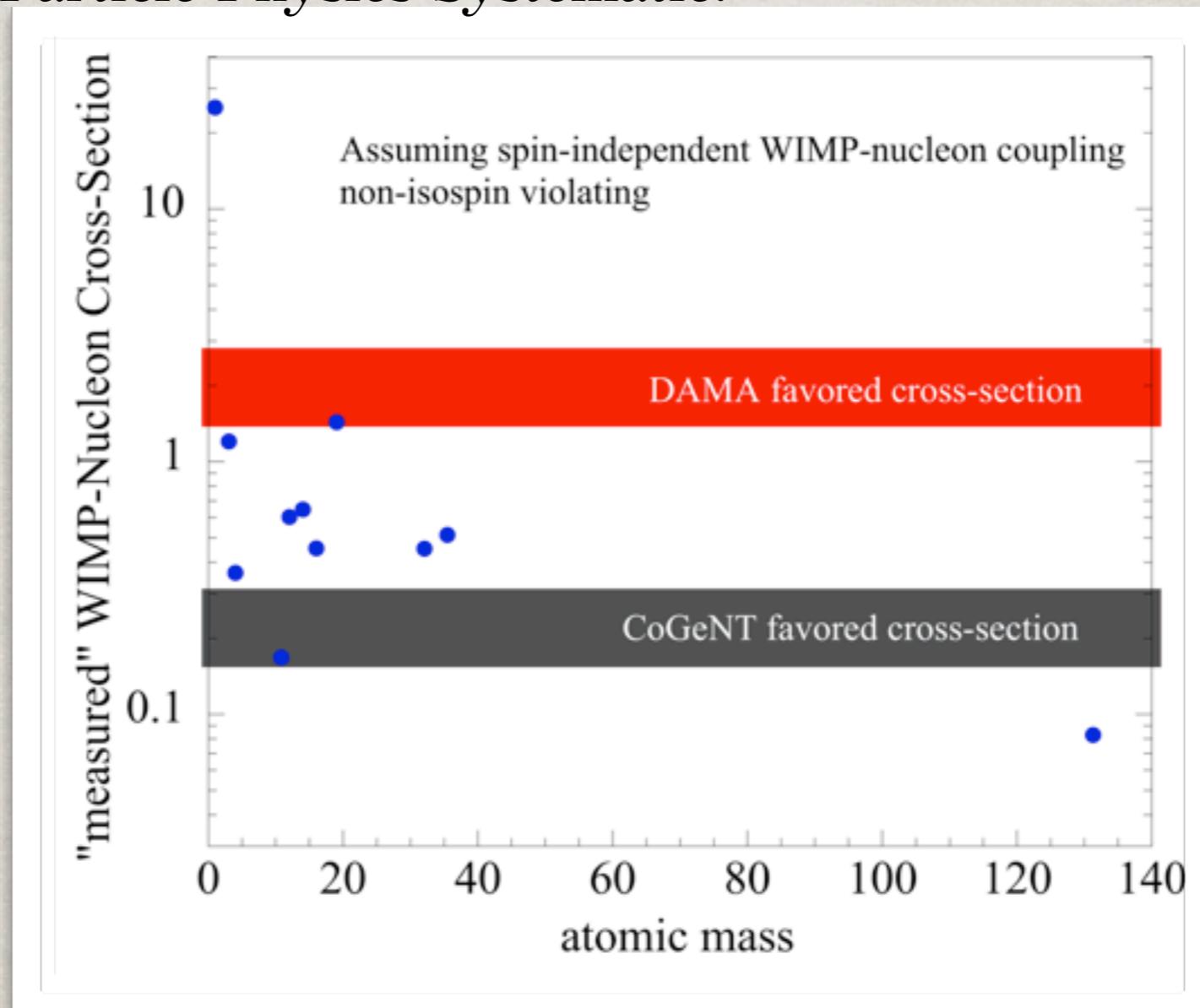
Light WIMPS

Study cross-section versus target mass.

- Search for (neutron) background systematic
- Characteristic coherence signal

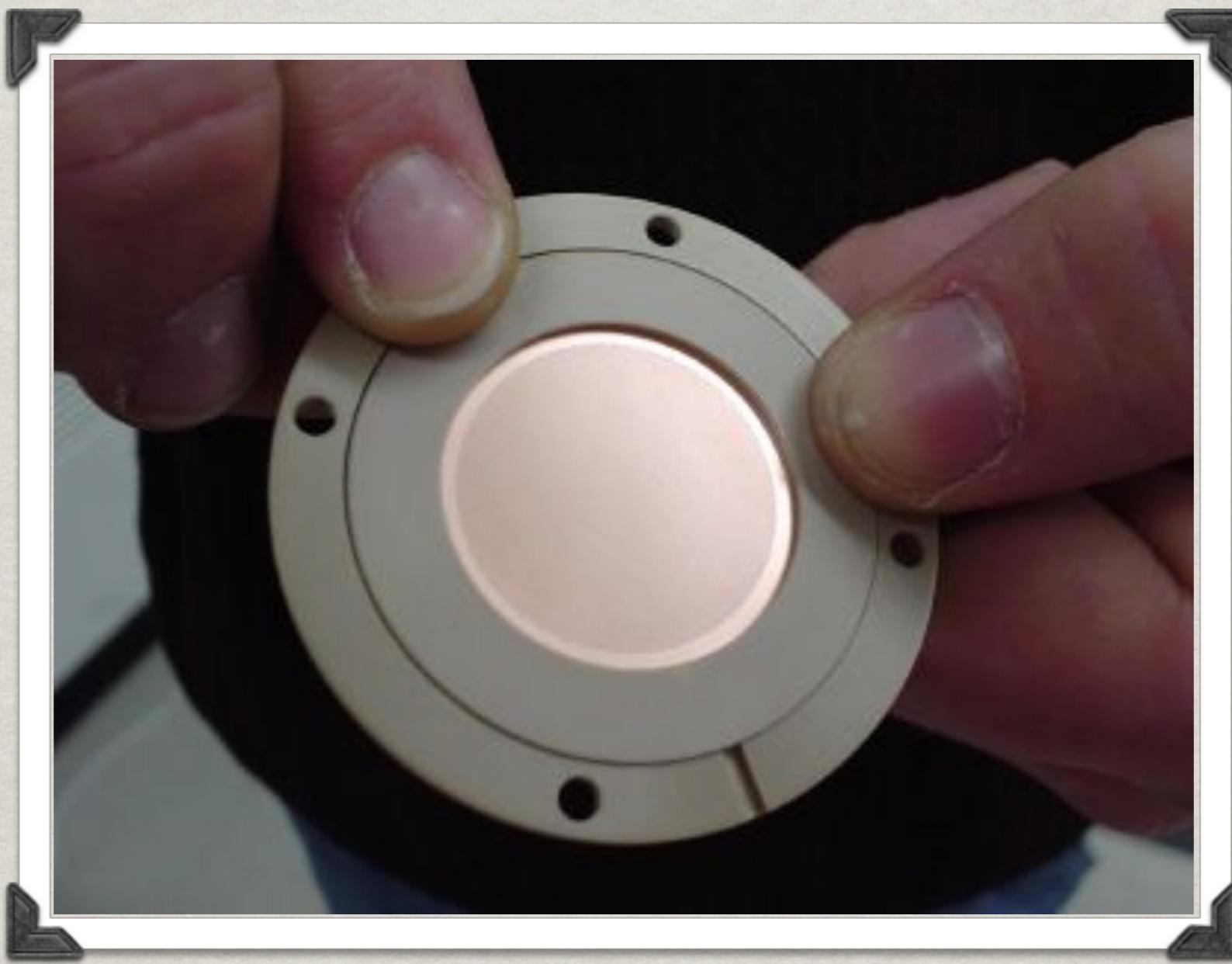
Cross-section cancelations can occur if we have isospin-violating WIMP interactions.

- Factorize out Particle Physics Systematic.

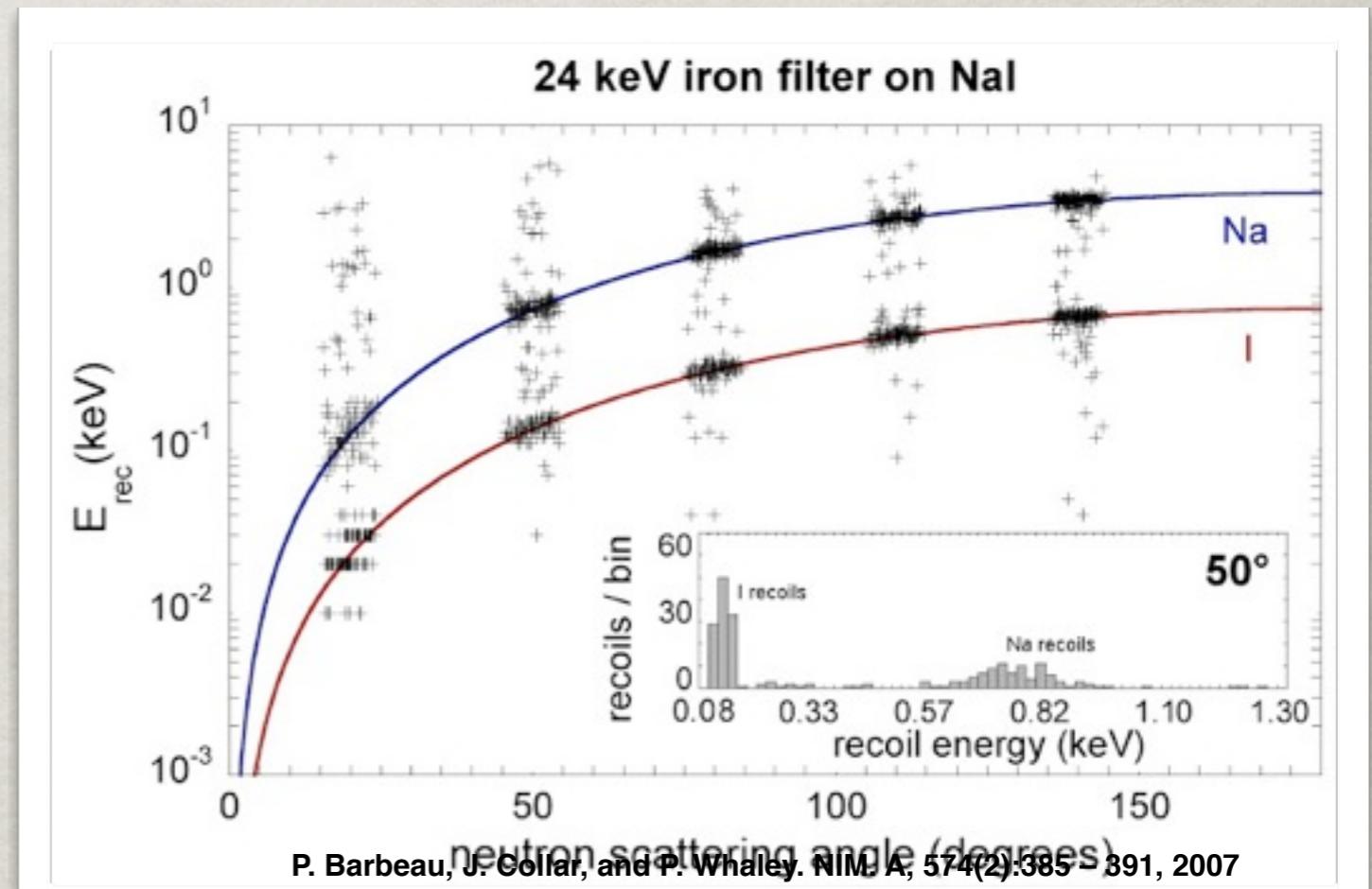
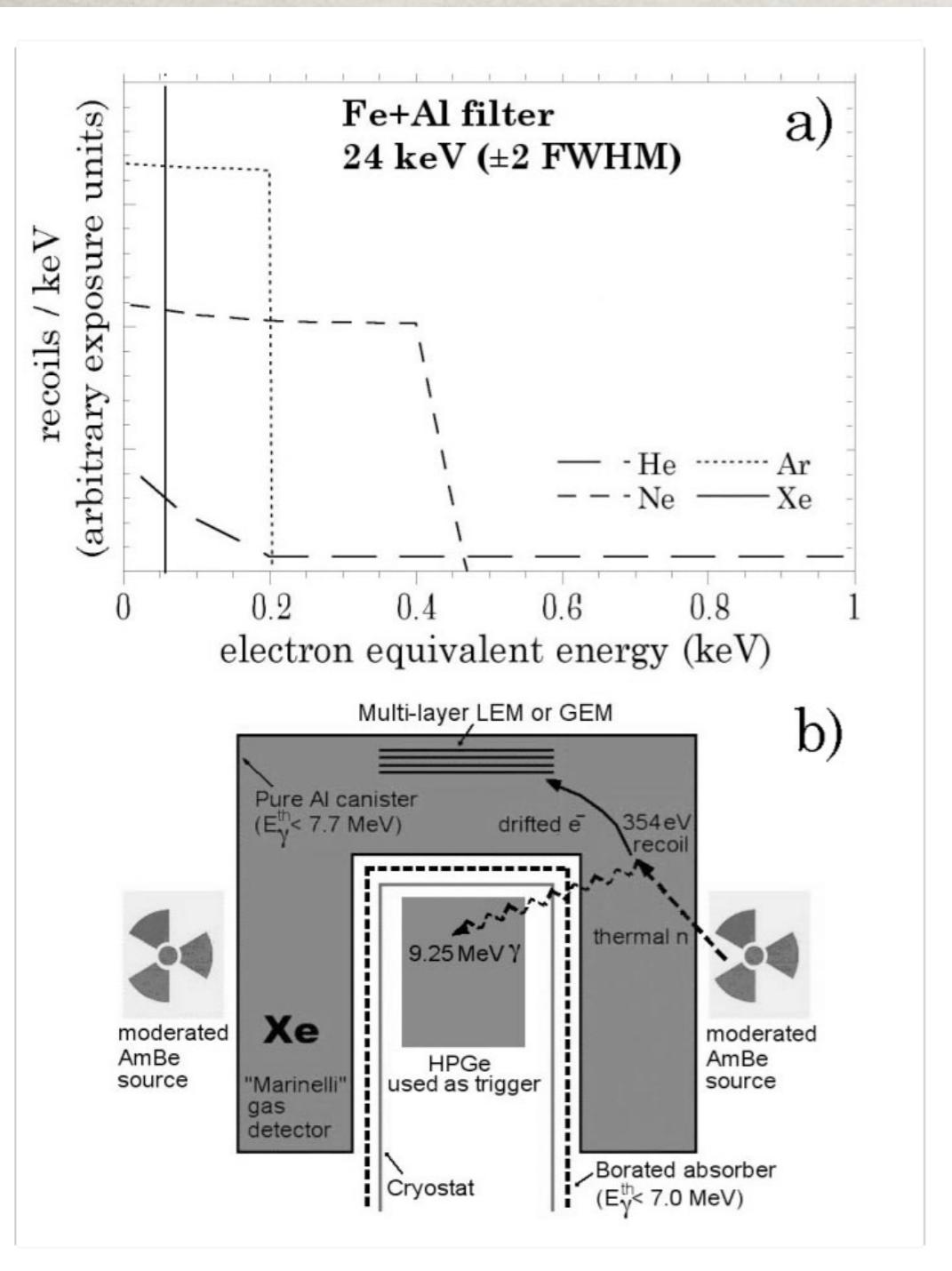


SUMMARY & OPEN ISSUES

- ✿ A detector concept has been presented which focuses on eliminating systematics with a simple/robust technology for precision CNNS & WIMP experiments
- ✿ Can we really predict the relative QF between ^{20}Ne and ^{22}Ne based on kinematics?
- ✿ Is there any ionization signal at all at low Q ?
- ✿ How difficult to enrich to ^{22}Ne ?
- ✿ High precision calibration of energy scale/electron gain...laser calibration?

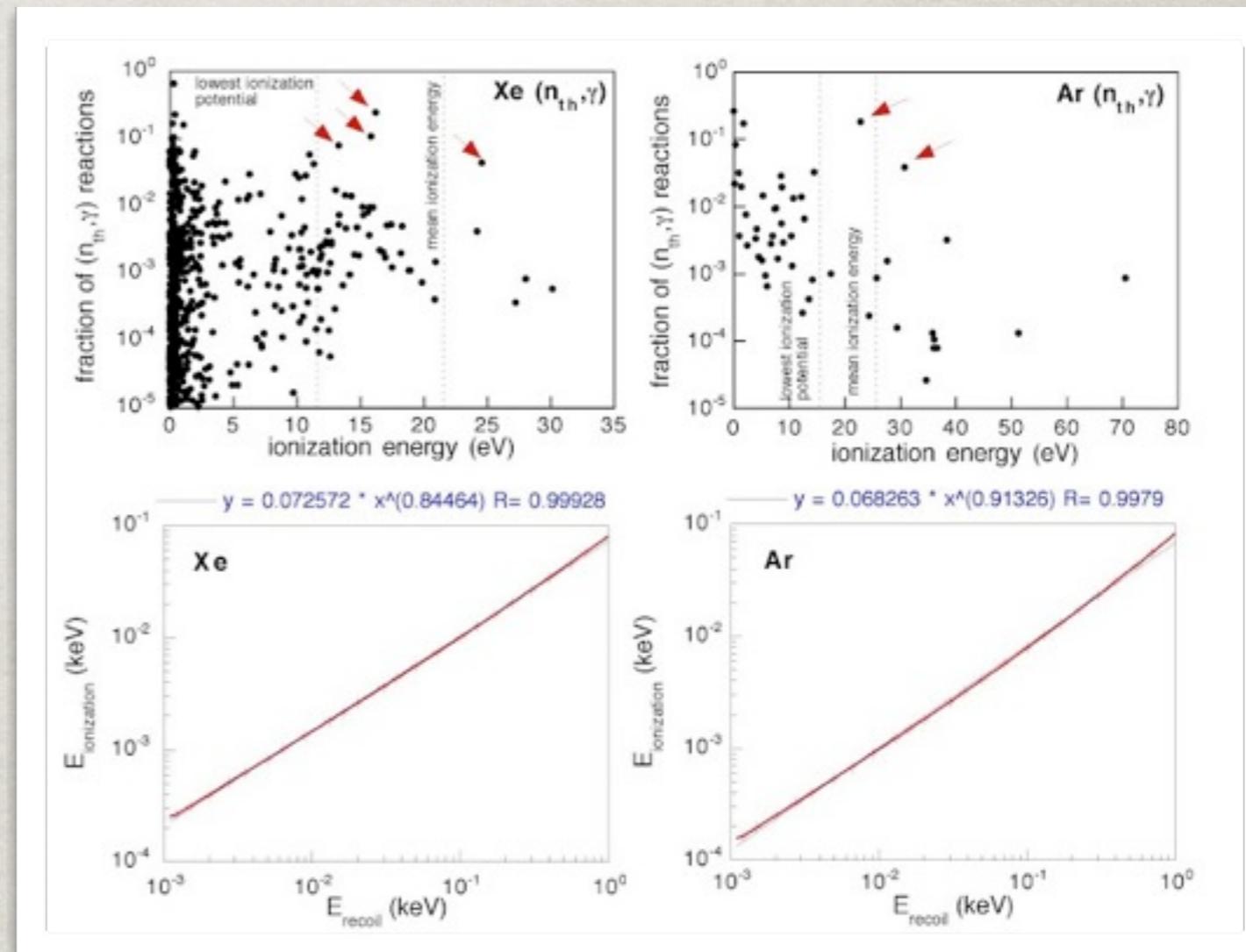
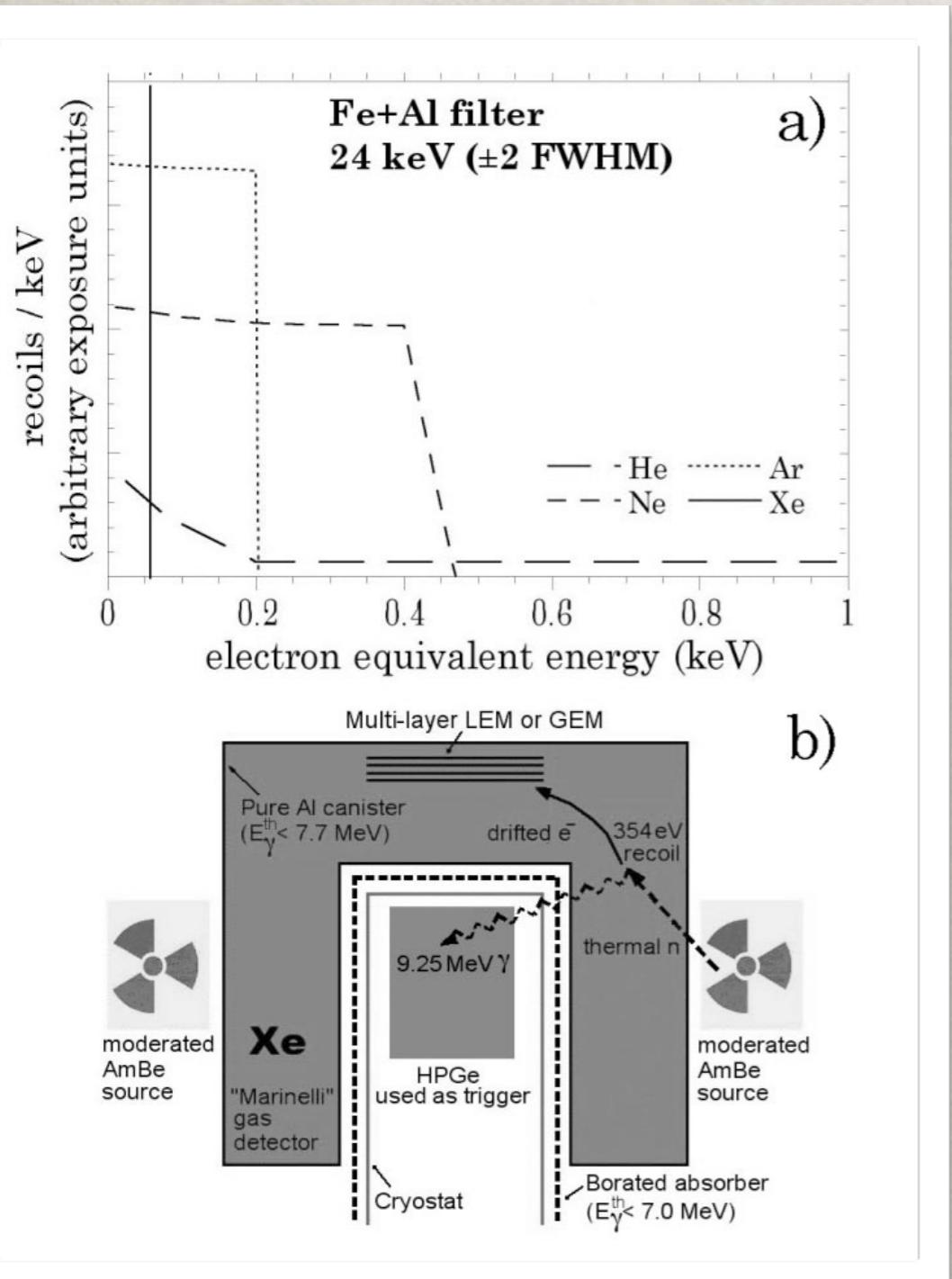


QF MEASUREMENT: 24 KEV MONOCHROMATIC NEUTRON BEAM (KSU)



P. Barbeau, J. I. Collar, J. Miyamoto, and I. Shipsey. IEEE Trans. Nucl. Sci., 50:1285–1289, 2003.

QF MEASUREMENT: THERMAL NEUTRONS



P. Barbeau, J. I. Collar, J. Miyamoto, and I. Shipsey. IEEE Trans. Nucl. Sci., 50:1285–1289, 2003.